

**NRC-CNRC**

Canadian Neutron  
Beam Centre



University of Connecticut

# *SMALL ANGLE NEUTRON SCATTERING (SANS)*

**Mu-Ping Nieh** (聂睦平)

*Institute of Materials Science &  
Dept. Chem. & Biomol. Eng.  
Univ. of Connecticut, Storrs, CT, USA*

*Presented to CNBC, 2013 Neutron Scattering Summer School, June 5*



National Research  
Council Canada

Conseil national  
de recherches Canada

Canada 

- Common Features of SANS
- Scattering Principle
- SANS Instrument & Data Reduction
- Contrast
- Dilute Particulate Systems – Form Factor
  - Data analysis
  - Examples
- Concentrated Particulate Systems – Structure Factor
  - Data analysis
  - Examples
- Dilute Polymer Solutions
  - Analysis for Gaussian Chain – Debye Function
  - More Generalized Analysis – Zimm Plot
- Small Angle Diffraction Analysis
- Model Independent Scaling Method Analysis
- Summary

# Common Features & Advantages of SANS

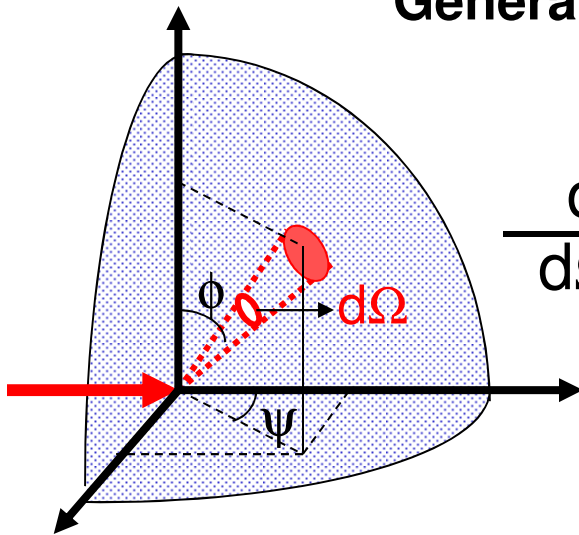
1. The  $Q$  range of SANS : between 0.002 and 0.6  $\text{\AA}^{-1}$   
– achievable with **long** wavelengths of neutrons and **low** detecting angles
2. A powerful technique for *in-situ* study on the *global* structures of *isotropic* samples
3. Easy to play contrast with isotope substitution
4. Data analysis:
  - (A) Model dependent: well-defined particulate systems  
(possible non-unique solutions)
  - (B) Model independent: general feature of the structure

# Neutron Scattering Principle - I

## General Equation

The number of neutrons scattered per second into a solid angle  $d\Omega$  with the final energy between  $E$  and  $E+dE$

$$\frac{d^2\sigma}{d\Omega dE} = \frac{\text{neutron flux of the incident beam} \cdot d\Omega dE}{\text{neutron flux of the incident beam} \cdot d\Omega dE}$$



## For Elastic Scattering

We do not analyze the energy but only count the number of the scattered neutrons

$$\frac{d\sigma}{d\Omega} = \int_0^{\infty} \frac{d^2\sigma}{d\Omega dE} dE = \frac{\text{The number of neutrons scattered per second into a solid angle } d\Omega}{\text{neutron flux of the incident beam} \cdot d\Omega}$$

[time<sup>-1</sup>]

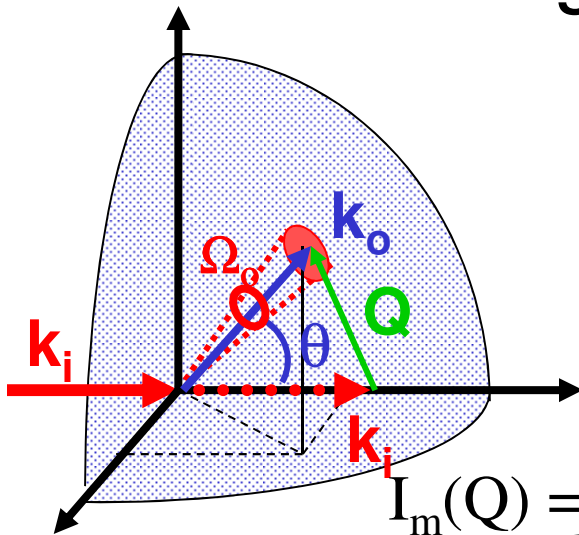
[time<sup>-1</sup>area<sup>-1</sup>] [area]

Differential cross section

# Neutron Scattering Principle - II

Scattering vector,  $Q$

$$|\mathbf{Q}| = |\mathbf{k}_o - \mathbf{k}_i| = \frac{4\pi}{\lambda} \sin \frac{\theta}{2}$$



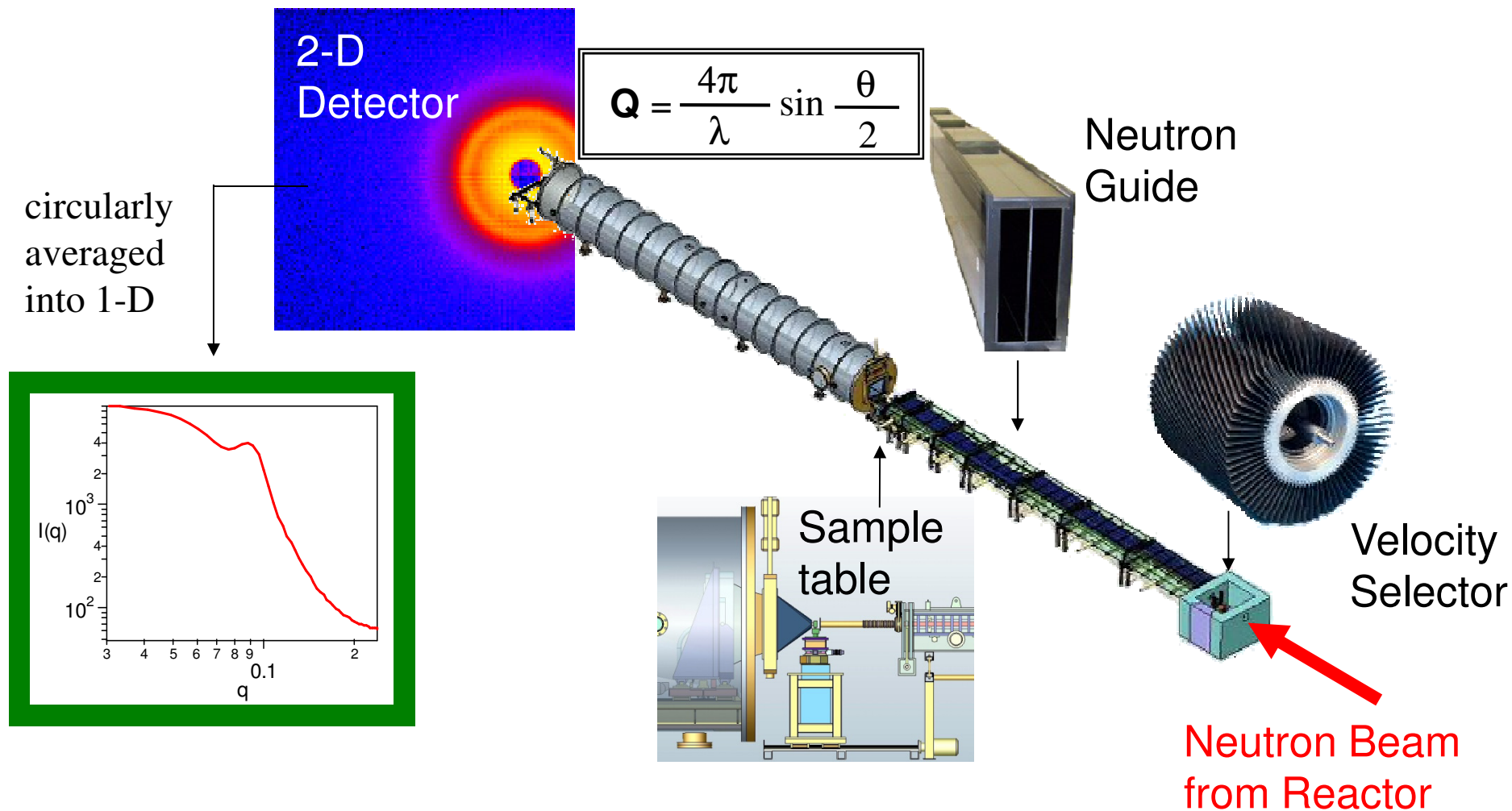
The measured intensity at  $Q$  (or  $\theta$ ),  $I_m(Q)$  can be expressed as

$$I_m(Q) = I_F \cdot \Omega_o \cdot \varepsilon \cdot T \cdot \left( \frac{d\sigma}{d\Omega} \right)_v \cdot A \cdot t$$

Flux  $\rightarrow I_F$     Solid angle  $\rightarrow \Omega_o$     Detector efficiency  $\rightarrow \varepsilon$     Sample transmission  $\rightarrow T$     Differential cross-section  $\rightarrow \left( \frac{d\sigma}{d\Omega} \right)_v$     Beam area on the sample  $\rightarrow A$     Path length  $\rightarrow t$

$$\left( \frac{d\sigma}{d\Omega} \right)_{v,sam} = \left( \frac{d\sigma}{d\Omega} \right)_{v,std} \frac{I_{m,sam}(Q) \cdot A_{std} \cdot t_{std} \cdot T_{std}}{I_{m,std}(Q) \cdot A_{sam} \cdot t_{sam} \cdot T_{sam}}$$

# Typical SANS Instrument



For  $\lambda = 8 \text{ \AA}$  and  $\theta_{\min} \sim 0.25^\circ$ ,  $Q_{\min} \sim 0.003 \text{ \AA}^{-1}$ .  
 Max. attainable length scale =  $2\pi/Q_{\min} \sim 2000 \text{ \AA}$ .

## Current CNBC SANS Setup

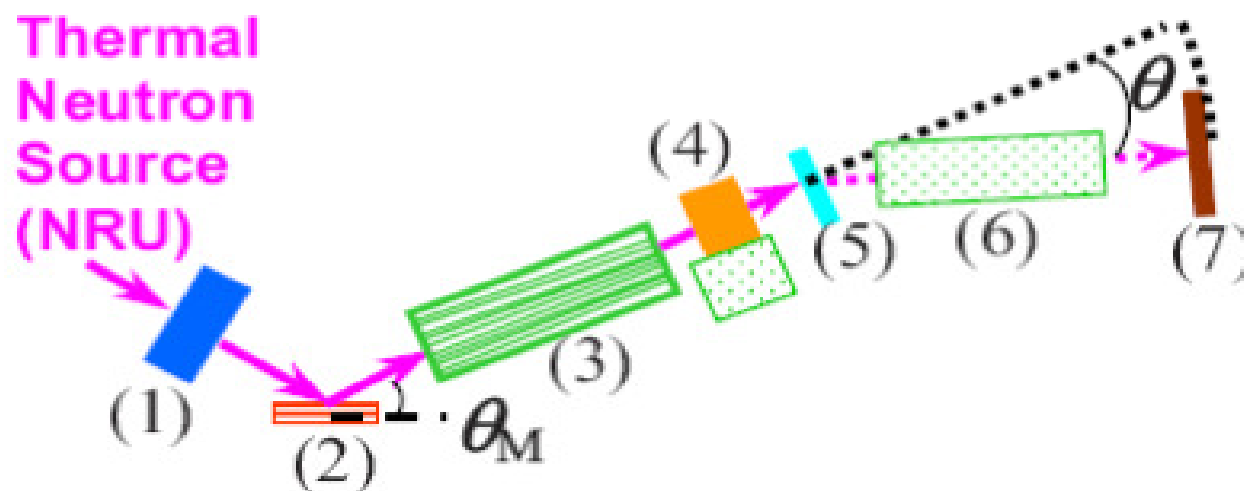


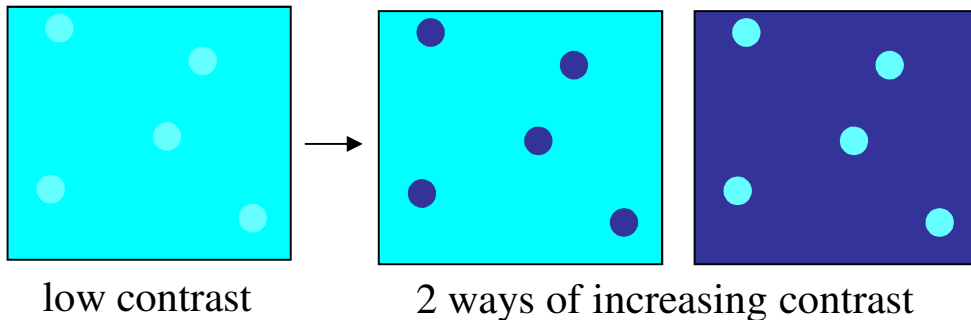
FIG. 1. (Color online) Schematic of the N5-SANS adapted from a triple-axis spectrometer to an instrument capable of SANS measurements. The components are as follows: (1) sapphire or Be filter. (2) Monochromator. (3) 23-channel CSC. (4) PG filter/21.6 cm long HSC/open. (5) Sample. (6) 48 cm long HSC. (7) 32-wire position sensitive detector.

M.-P. Nieh Y. Yamani, N. Kučerka and J. Katsaras  
*Rev. Sci. Instrum.* (2008) 79, 095102

$Q_{\min} \sim 0.006 \text{ \AA}^{-1}$ : Max. attainable length =  $2\pi/Q_{\min} \sim 1000 \text{ \AA}$ .

# Contrast – A Key Parameter

$$\left(\frac{d\sigma}{d\Omega}\right)_V = \frac{N}{V} (\rho_p - \rho_o)^2 V_p^2 P(\vec{Q}) S(\vec{Q})$$

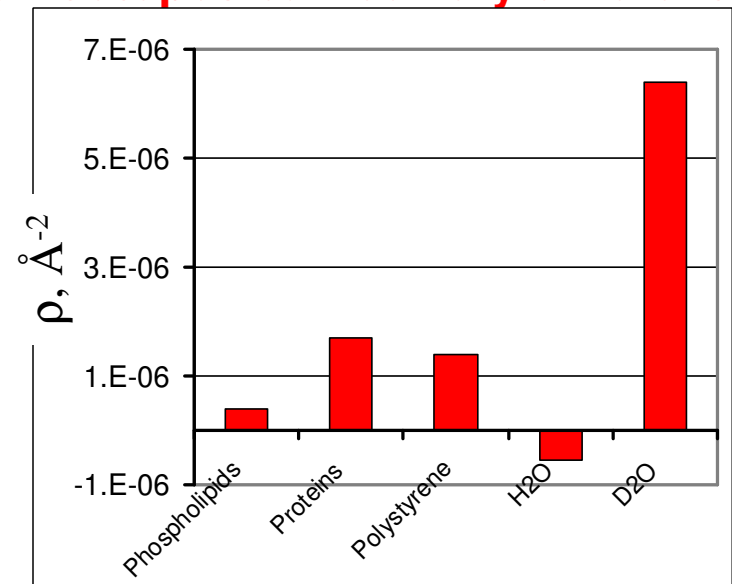


Scattering intensity is proportional to “the square of the scattering length density difference between studied materials and medium” (known as “contrast factor”).

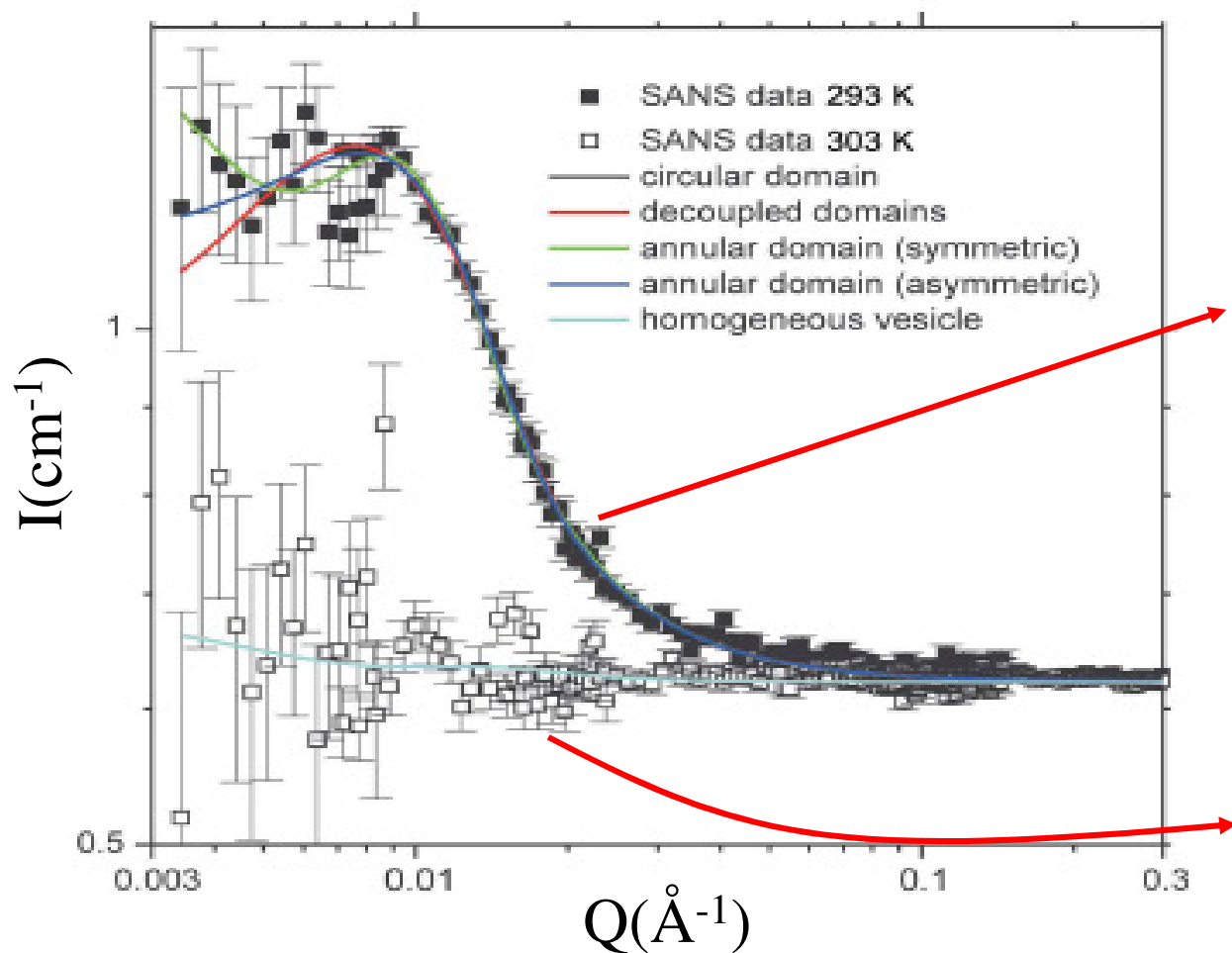
**Ideally, we would like to increase the contrast yet without changing the chemical properties of the system. This is one of the greatest advantages using neutron scattering, since the neutron scattering lengths of isotopes can be very different.**

## Biomaterials

Atom	Nucleus	$b_c$ ( $10^{-12}$ cm)	$\sigma_c$ ( $10^{-24}$ cm <sup>2</sup> )	$\sigma_l$ ( $10^{-24}$ cm <sup>2</sup> )	$\sigma_{abs}^*$ ( $10^{-24}$ cm <sup>2</sup> )
Hydrogen	<sup>1</sup> H	-0.374	1.76	79.7	0.33
Deuterium	<sup>2</sup> H	0.667	5.59	2.01	0
Carbon	<sup>12</sup> C	0.665	5.56	0	0
Nitrogen	<sup>14</sup> N	0.930	11.1	0	1.88
Oxygen	<sup>16</sup> O	0.580	4.23	0	0
Fluorine	<sup>19</sup> F	0.556	4.03	0	0
Phosphorous	P <sup>†</sup>	0.513	3.31	0	0.17
Chlorine	Cl <sup>†</sup>	0.958	11.53	5.9	33.6







J. Pencer, V. N. P. Anghel, N. Kučerka and J. Katsaras,  
*J. Appl Cryst* 40 (2007) 513-525

$$\left(\frac{d\sigma}{d\Omega}\right)_v = \frac{N}{V} (\rho_p - \rho_o)^2 V_p^2 P(\vec{Q}) S(\vec{Q})$$

orientational average

$$\left(\frac{d\sigma}{d\Omega}\right)_v = \frac{N}{V} \iint \rho(\vec{r}_1)\rho(\vec{r}_2) \langle e^{-i\vec{Q}\cdot(\vec{r}_1-\vec{r}_2)} \rangle d\vec{r}_1 d\vec{r}_2$$

$$= \frac{N}{V} \left| \int_{V_p} \Delta\rho(\vec{r}) \cdot \langle e^{-i\vec{Q}\cdot\vec{r}} \rangle d\vec{r} \right|^2$$

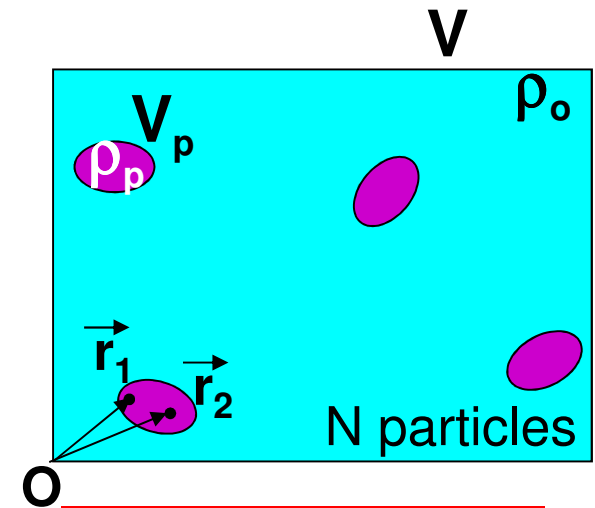
$$= \frac{N}{V} (\rho_p - \rho_o)^2 \cdot V_p^2 \cdot P(\vec{Q})$$

$$= \frac{(\rho_p - \rho_o)^2 N}{V} \sum_{k=1}^{2N} |A_k(\vec{Q})|^2$$

Particle Form Factor

Amplitude of the Form Factor

$$A_k(\vec{Q}) = \int_{\text{particle } k} \langle e^{-i\vec{Q}\cdot\vec{r}} \rangle d\vec{r}$$



$\vec{r}_1$  and  $\vec{r}_2$  are vectors of scattering center from one particle

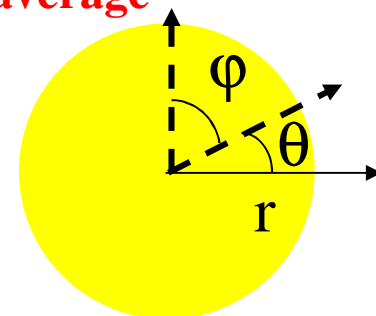
The form factor is determined by the structure of the particle.

## Simulation of a particular system - Spheres

$$P(q) = \frac{1}{V_{\text{sphere}}^2} \left| \iiint_{V_{\text{sphere}}} e^{-i\vec{Q} \cdot \vec{r}} d\vec{r} \right|^2$$

Since spheres are isotropic, there is no need to do orientational average

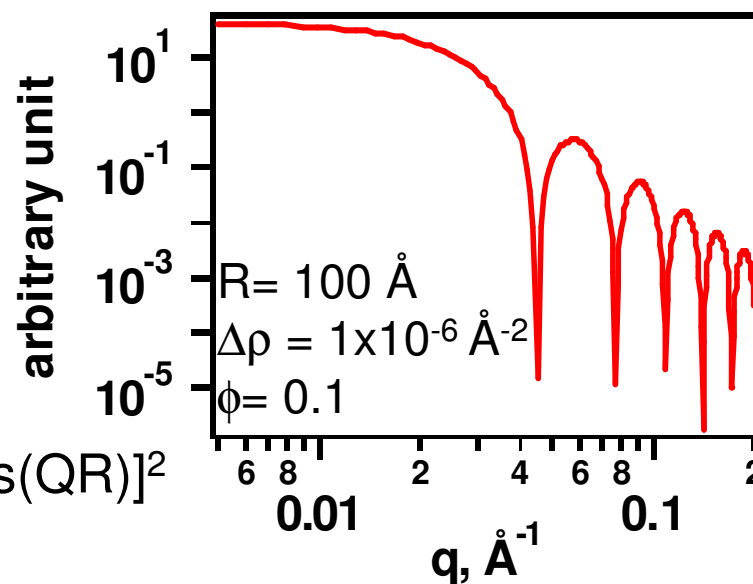
$$= \frac{1}{V_{\text{sphere}}^2} \left| \int_{r=0}^{r=R} \int_{\theta=0}^{\theta=\pi} \int_{\phi=0}^{\phi=2\pi} e^{-iQr \cdot \cos\theta} r^2 \sin\theta d\phi d\theta dr \right|^2$$




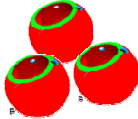
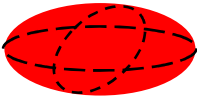
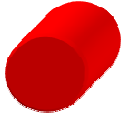


$$= \frac{9}{(QR)^6} [\sin(QR) - QR \cdot \cos(QR)]^2$$

$$\left( \frac{d\sigma}{d\Omega} \right)_v = \frac{N}{V} (\rho_{\text{sphere}} - \rho_o)^2 \cdot V_{\text{sphere}}^2 \cdot P(\vec{Q})$$

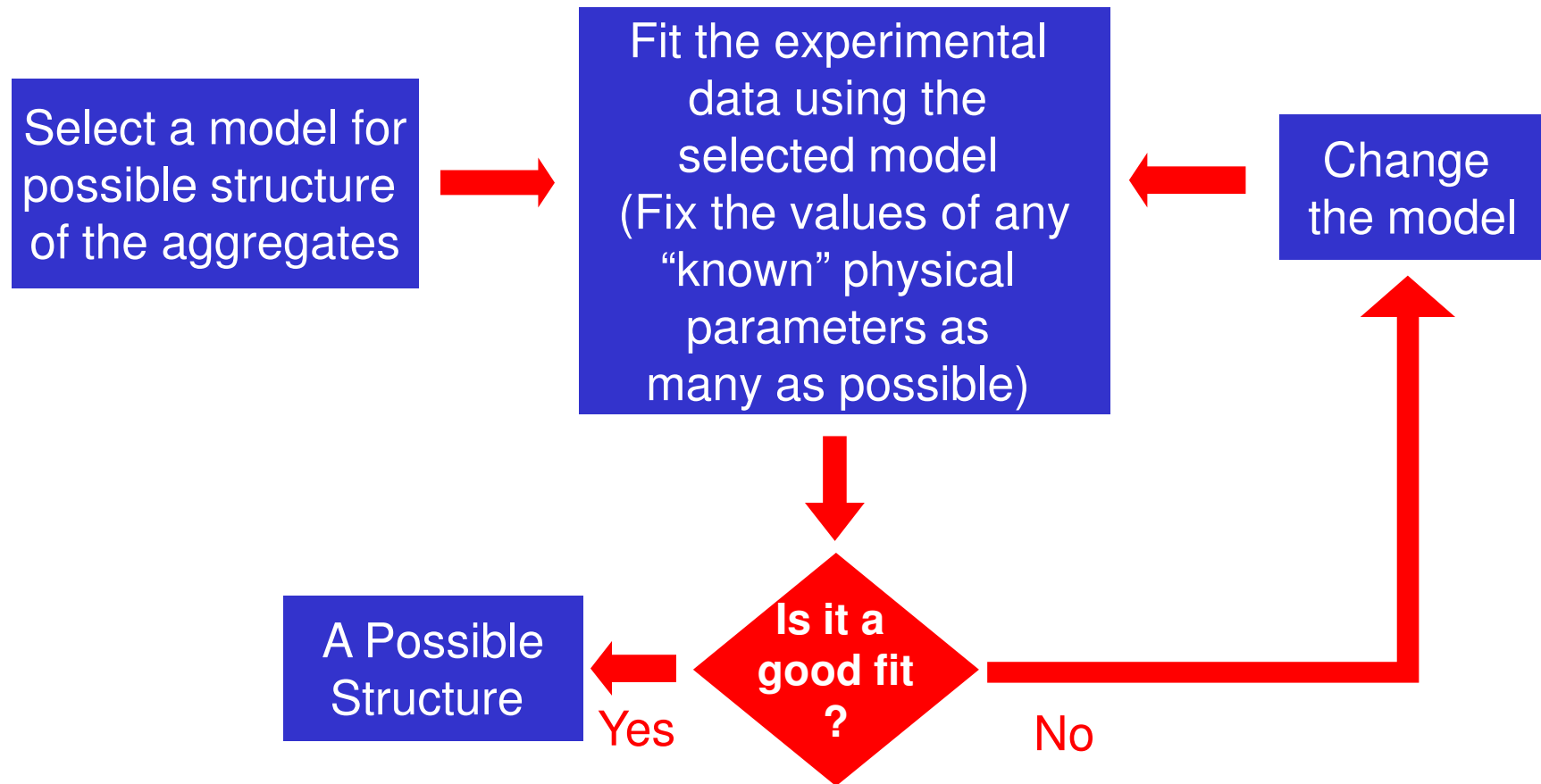
$$= \phi_{\text{sphere}} V_{\text{sphere}} \Delta\rho^2 \frac{9}{(QR)^6} [\sin(QR) - QR \cdot \cos(QR)]^2$$



## Common form factors of particulate systems

Morphologies	P(Q)	Morphologies
<b>Spheres</b> (radius :R)	$\frac{9}{(QR)^6} [\sin(QR) - QR \cdot \cos(QR)]^2 = A_{\text{sph}}^2(QR)$	
<b>Spherical shells</b> (outer radius: R <sub>1</sub> inner radius: R <sub>2</sub> )	$\frac{[R_1^3 \cdot A_{\text{sph}}(QR_1) - R_2^3 \cdot A_{\text{sph}}(QR_2)]^2}{(R_1^3 - R_2^3)^2}$	
<b>Triaxial ellipsoids</b> (semiaxes: a,b,c)	$\int_0^1 \int_0^1 A_{\text{sph}}^2 [Q \sqrt{a^2 \cos^2(\pi x/2) + b^2 \sin^2(\pi x/2)(1-y^2)_1 + c^2 y^2}] dx dy$	
<b>Cylinders</b> (radius: R length: L)	$4 \int_0^1 \frac{J_1^2[QR\sqrt{1-x^2}]}{[QR\sqrt{1-x^2}]^2} \frac{\sin^2(QLx/2)}{(QLx/2)^2} dx$ <i>J<sub>1</sub>(x) is the first kind Bessel function of order 1</i>	
<b>Thin disk</b> (radius: R)	By setting L = 0 $\frac{2 - J_1(2QR)/QR}{Q^2 R^2}$	
<b>Long rod</b> (length: L)	By setting R = 0 $\frac{2}{QL} \int_0^{QL} \frac{\sin(t)}{t} dt - \frac{\sin^2(QL/2)}{(QL/2)^2}$	

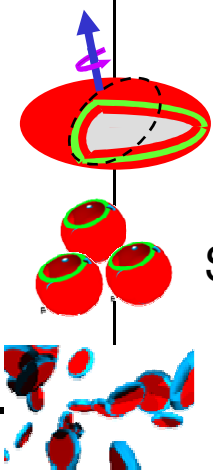
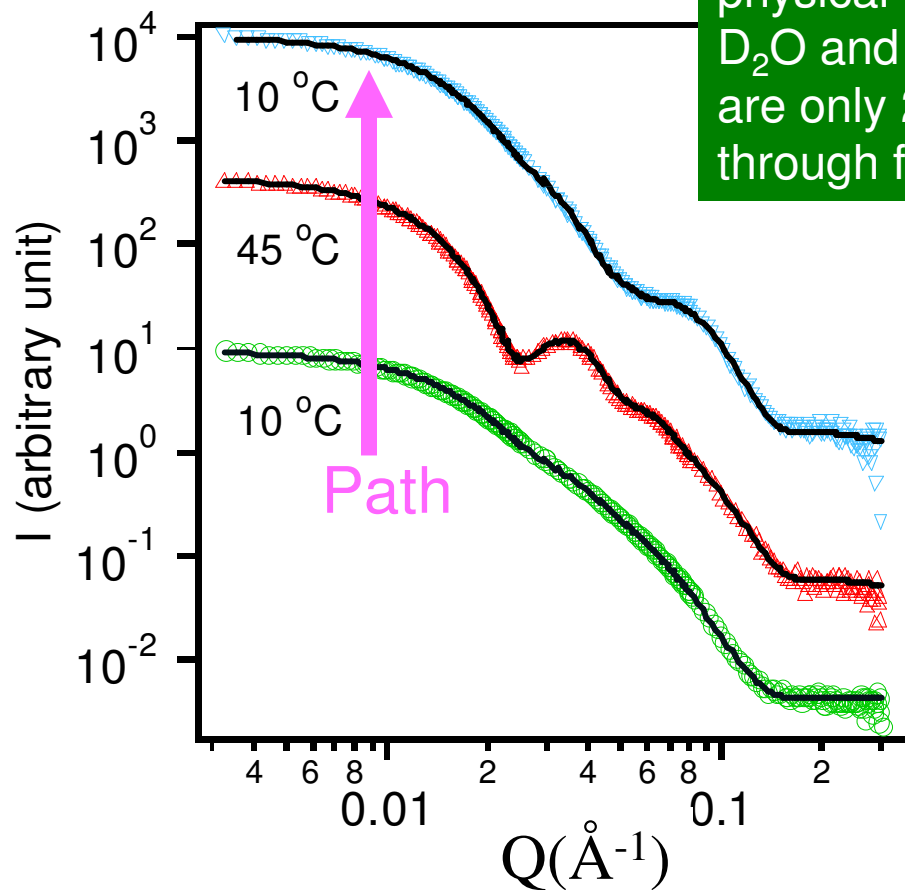
## *Procedure of Data Analysis*



## Examples of A Dilute Particulate System

Sample: A mixture of dimyristoyl and dihexanoyl phosphatidylcholine (DMPC, DHPC), and dimyristoyl phosphatidylglycerol (DMPG) in D<sub>2</sub>O;  
total lipid conc. = 0.1 wt.%

In these cases, the known (or constrained) physical parameters are SLDs of lipid and D<sub>2</sub>O and the bilayer thickness. Thus, there are only 2 parameters to be determined through fitting in each case.



Oblate shells:  $a=180 \text{ \AA}$ ,  $b=62 \text{ \AA}$

Spherical shells:  $R=133 \text{ \AA}$ ,  $p= 0.15$

Bilayer disks:  $R=156 \text{ \AA}$ ,  $L= 45 \text{ \AA}$

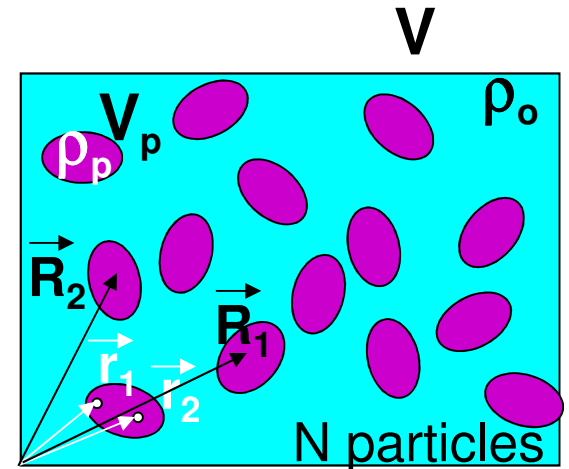
# SANS Analysis – Concentrated Particulate Systems

$$\left( \frac{d\sigma}{d\Omega} \right)_V = \frac{N}{V} (\rho_p - \rho_o)^2 V_p^2 P(\vec{Q}) S(\vec{Q})$$

$$\left( \frac{d\sigma}{d\Omega} \right)_V = \frac{N}{V} \iint \rho(\vec{r}_1) \rho(\vec{r}_2) \langle e^{-i\vec{Q} \cdot (\vec{r}_1 - \vec{r}_2)} \rangle d\vec{r}_1 d\vec{r}_2$$

$$= \frac{1}{V} \sum_{k=1}^N \langle |A_k(\vec{Q})|^2 \rangle$$

$$+ \frac{1}{V} \left\langle \sum_{k=1}^N \sum_{j \neq k}^N A_k(\vec{Q}) A_j^*(\vec{Q}) e^{-i\vec{Q} \cdot (\vec{R}_k - \vec{R}_j)} \right\rangle$$



$$\sim \frac{1}{V} \sum_{k=1}^N \langle |A_k(\vec{Q})|^2 \rangle \left\{ 1 + \sum_{j \neq k}^N \langle e^{-i\vec{Q} \cdot (\vec{R}_k - \vec{R}_j)} \rangle \right\}$$

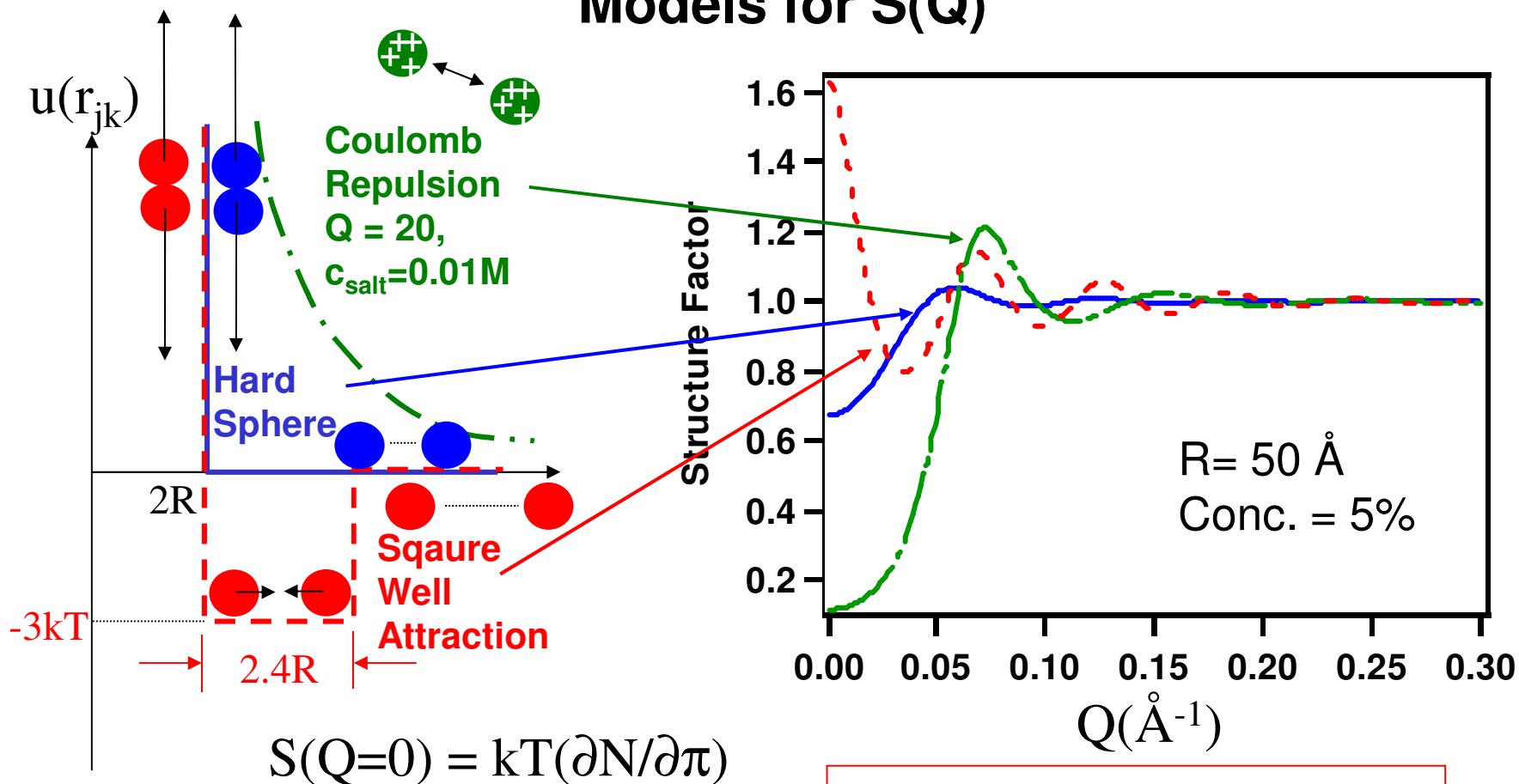
$$= \frac{N}{V} (\rho_p - \rho_o)^2 \cdot V_p^2 \cdot P(\vec{Q}) S(\vec{Q})$$

Structure Factor

In dilute solution,  $S(\vec{Q}) = 1$

# SANS Analysis – Concentrated Particulate Systems

## Models for $S(Q)$



$$S(Q=0) = kT(\partial N/\partial \pi)$$

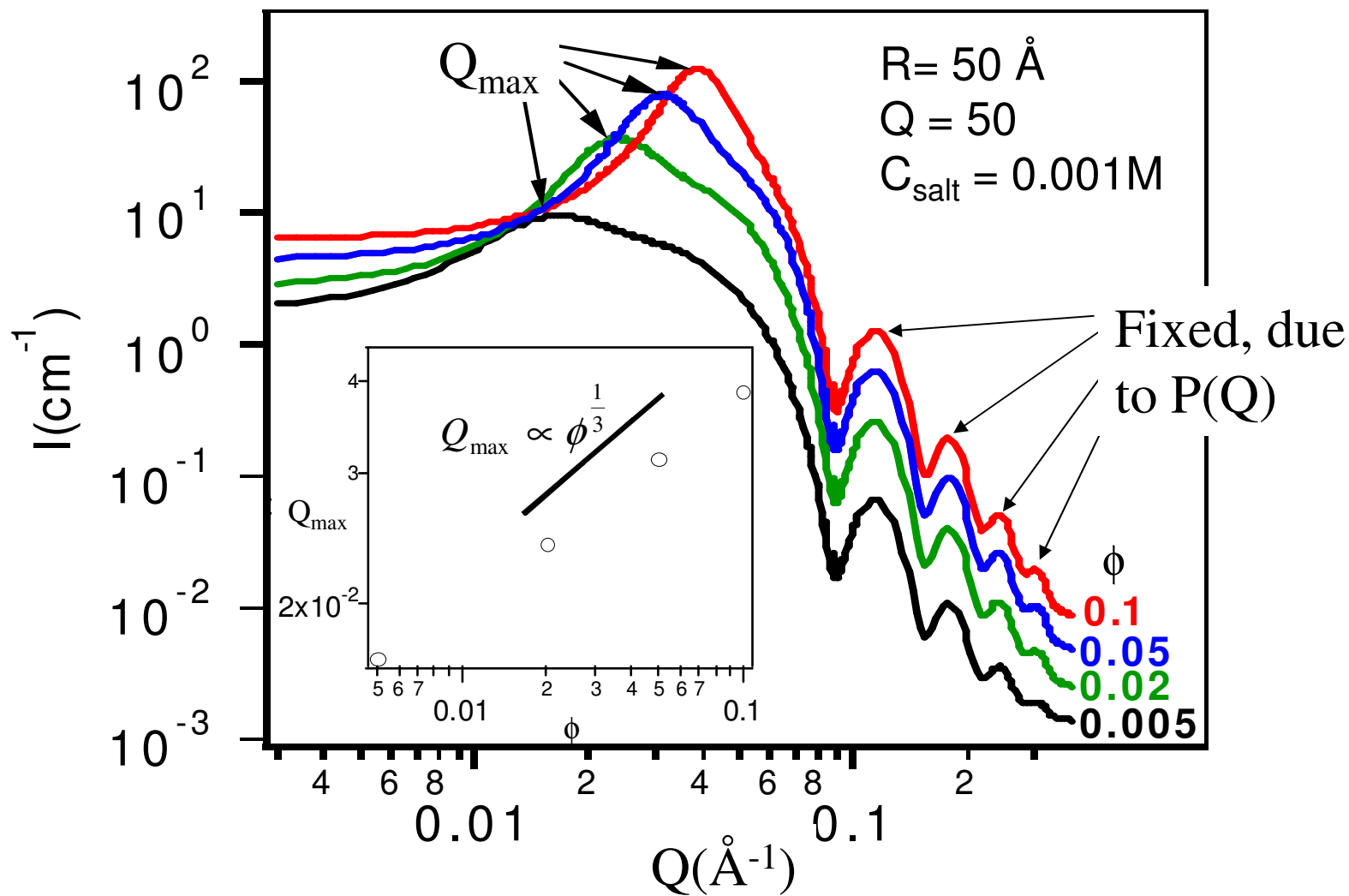
Osmotic compressibility  $\rightarrow$

$S(0) > 1$  more compressible  
 $S(0) < 1$  less compressible



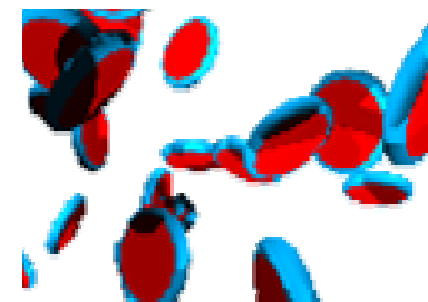
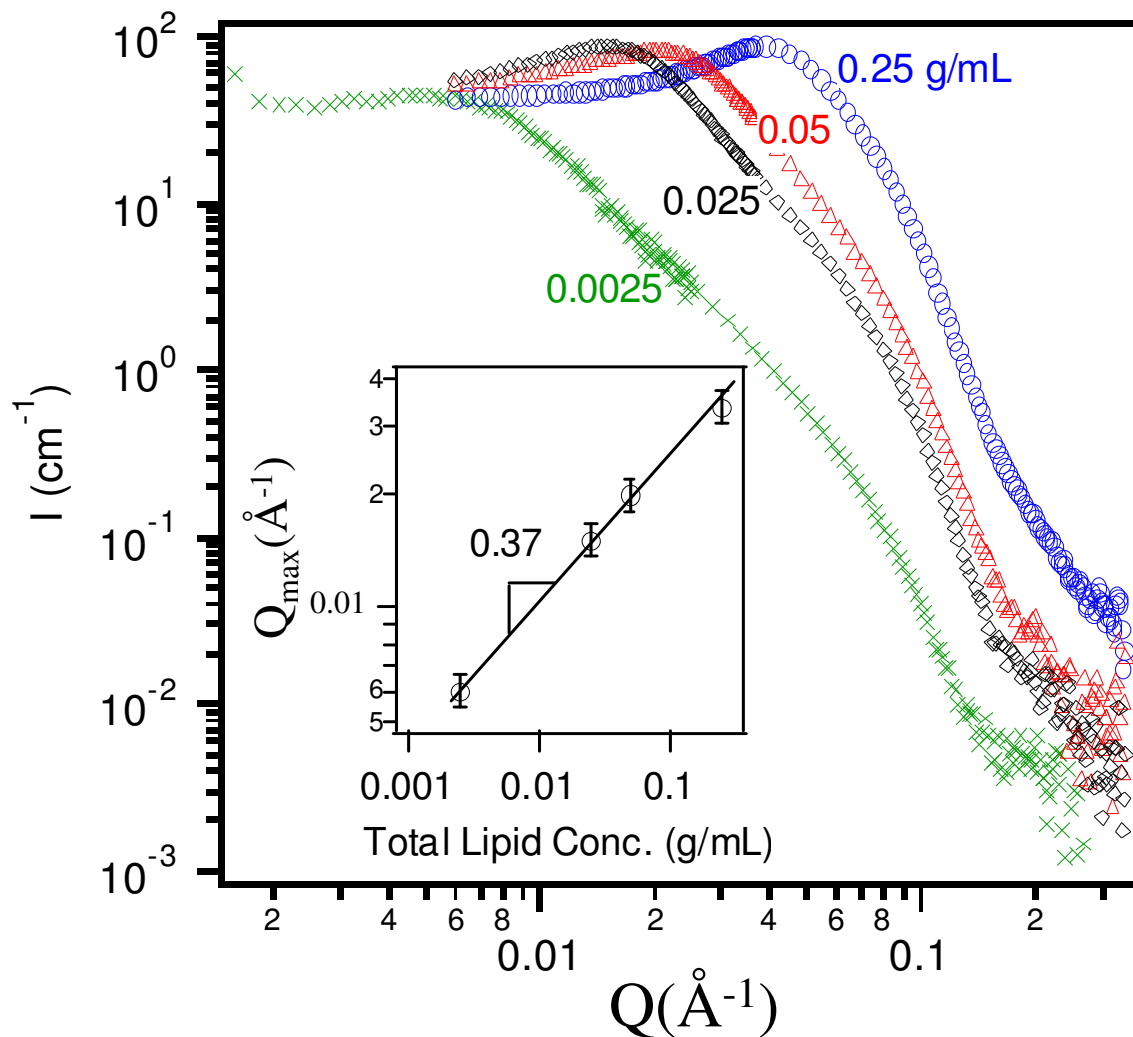
# SANS Analysis – Concentrated Particulate Systems

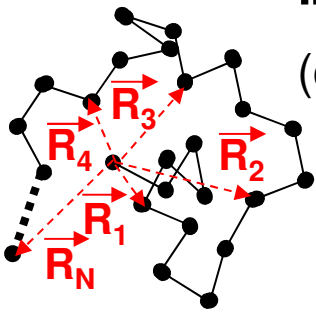
## Simulation Result using Coulomb Repulsive Model



# SANS Analysis – Concentrated Particulate Systems

“Bicelles” composed of DMPC/DHPC and small amount of  $Tm^{3+}$





**“Gaussian Chain” Model** : The effective bond length,  $a$ , of one step (composed of several segments of the chain) has a Gaussian distribution.

$$a^2 = \sum_{n=1}^N \frac{|\vec{R}_n - \vec{R}_{n-1}|^2}{N}$$

The distribution vector between any two “steps”  $m, n$ , i.e.,  $(\vec{R}_n - \vec{R}_m)$  is Gaussian

$$\Phi(\vec{R}_n - \vec{R}_m, n-m) = \left[ \frac{3}{2\pi a^2 |n-m|} \right]^{3/2} \exp\left[ -\frac{3(\vec{R}_n - \vec{R}_m)^2}{2a^2 |n-m|} \right]$$

This model is adequate for describing long polymer chains in a theta solvent, where the segment-segment and segment-solvent interaction and the excluded volume effect are canceled out. Such polymer chains are also known as “phantom” chains.

“The Theory of Polymer Dynamics” M. Doi and S. F. Edwards (1992)

# SANS Analysis – Dilute “Gaussian Chain” Solutions

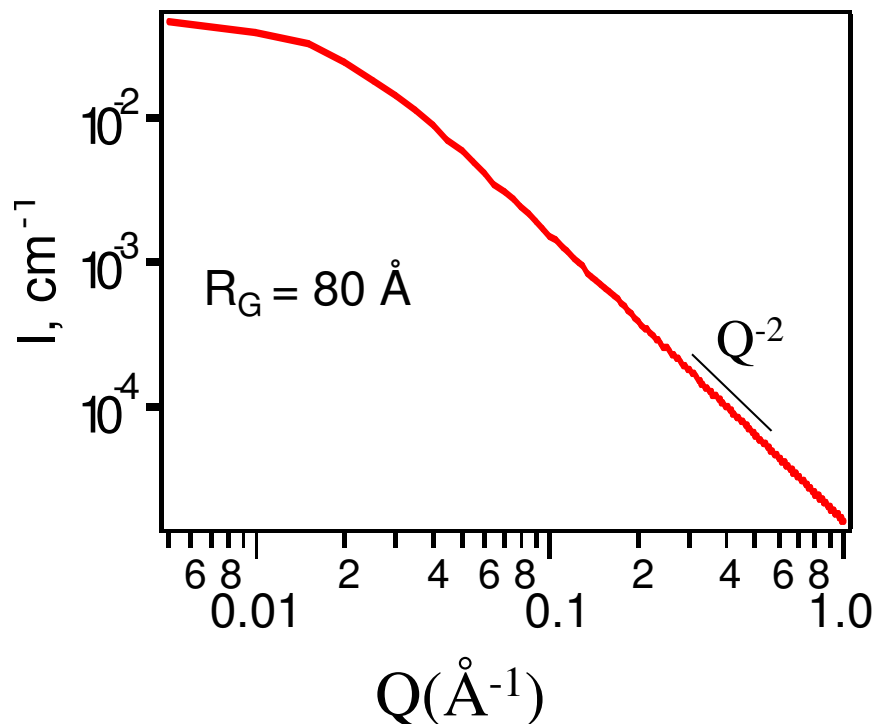
The scattering from a Gaussian Chain can be described as a Debye function

$$\left(\frac{d\sigma}{d\Omega}\right)_v = N_p \phi_p (\rho_p - \rho_o)^2 V_M \frac{2}{Q^4 R_G^4} (e^{-Q^2 R_G^2} - 1 + Q^2 R_G^2)$$

Debye Function

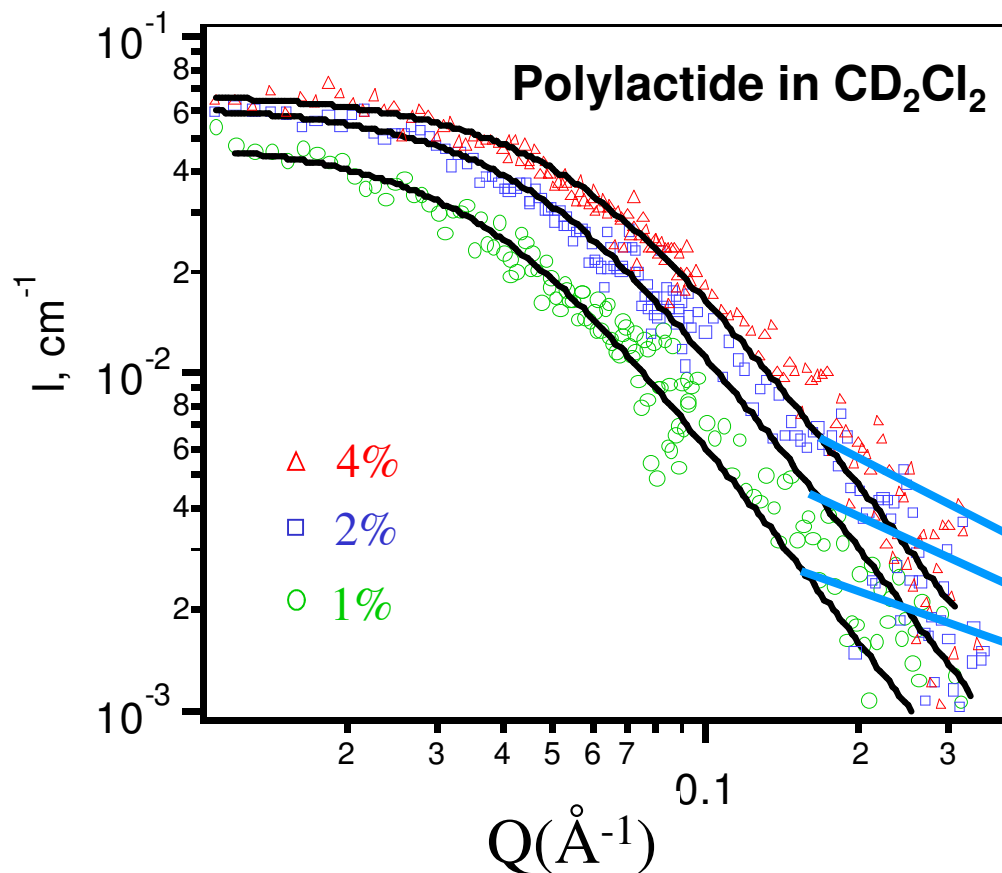
$V_M$ : molecular volume of monomer  
 $N_p$ : degree of polymerization

Radius of Gyration,  $R_G^2 = \frac{1}{2N^2} \sum_{m,n} \langle (R_n - R_m)^2 \rangle$



Debye function is not always adequate to describe the polymer conformation!

An example is that as polymers are in a good solvent, the intensity at high Q regime will decay with  $Q^{-5/3}$  instead of  $Q^{-2}$ .



$\phi_p$	$N_p$	$R_G$ (Å)
1%	111	39
2%	74	32
4%	42	27

**deviation from the data!**

For dilute polymer solutions, the scattering intensity at low  $q$  regime (i.e.,  $q \cdot R_G < 1$ ) can be expressed as a function of  $R_G$ ,  $M_W$  (molecular weight),  $A_2$  (second virial coefficient) and  $\phi_p$  (volume fraction of the polymer)

$$\frac{V_m \phi_p \Delta \rho^2}{I(Q)} = (1 + Q^2 R_G^2 / 3 + \dots) \left( \frac{1}{N_p} + 2 \phi_p A_2 + \dots \right)$$

$V_m$  is the molecular volume of the monomer  
 $N_p$  is the degree of polymerization

$R_G$ ,  $N_p$  and  $A_2$  can be obtained by Zimm plot, where

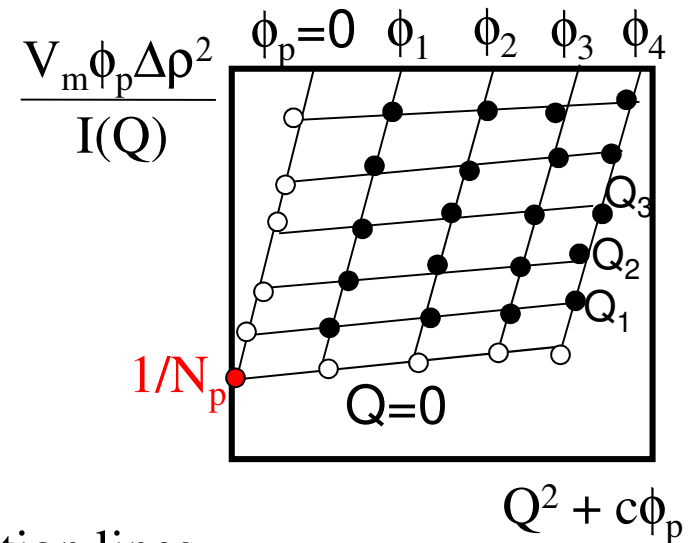
$\frac{V_m \phi_p \Delta \rho^2}{I(Q)}$  is plotted against  $(Q^2 + c \phi_p)$  and two extrapolation lines

for  $Q=0$  and  $\phi_p=0$  are also obtained ( $c$  is an arbitrary number.).

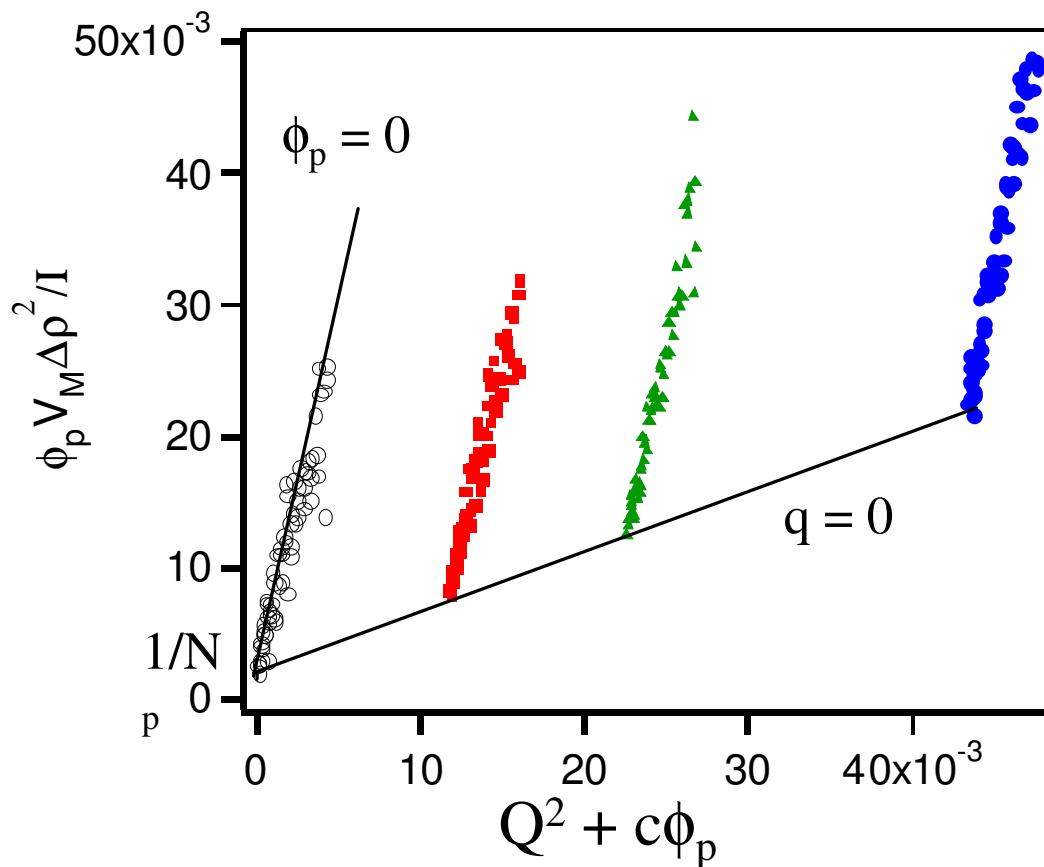
The slope of the extrapolation line for  $\phi_p = 0$  is  $R_G^2 / 3 N_p$ .

The slope of the extrapolation line for  $Q = 0$  is  $2 A_2 / c$ .

The intercept of both extrapolation lines is  $1 / N_p$ .



Poly lactide in  $CD_2Cl_2$



$N_p = 387$   
 $R_G = 79 \text{ \AA}$   
 $A_2 = 88.6$

Compared to the result obtained from Debye fitting

$\phi_p$	$N_p$	$R_G (\text{\AA})$
1%	111	39
2%	74	32
4%	42	27

It requires at least two concentrations to make up a Zimm plot.

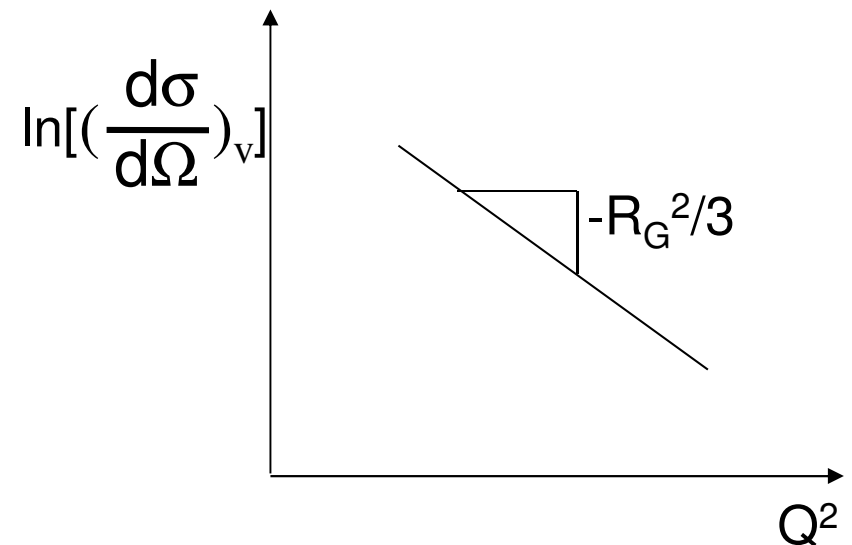
# SANS Analysis – Model Independent Scaling Method

*Guinier regime* – low-Q scattering

$$\left(\frac{d\sigma}{d\Omega}\right)_v = \frac{N}{V} (\rho_p - \rho_o)^2 \cdot V_p^2 \cdot [1 - (QR_G)^2/3 + \dots]$$

$$\sim \phi_p (\rho_p - \rho_o)^2 V_p e^{-Q^2 R_G^2/3}$$

$$\ln\left[\left(\frac{d\sigma}{d\Omega}\right)_v\right] = \ln(I_o) - Q^2 R_G^2/3$$



“Polymers and Neutron Scattering” J. S. Higgins and H. C. Benoit



## SANS Analysis – Model Independent Scaling Method

**Porod regime** – high-Q (with respect to the length scale) scattering

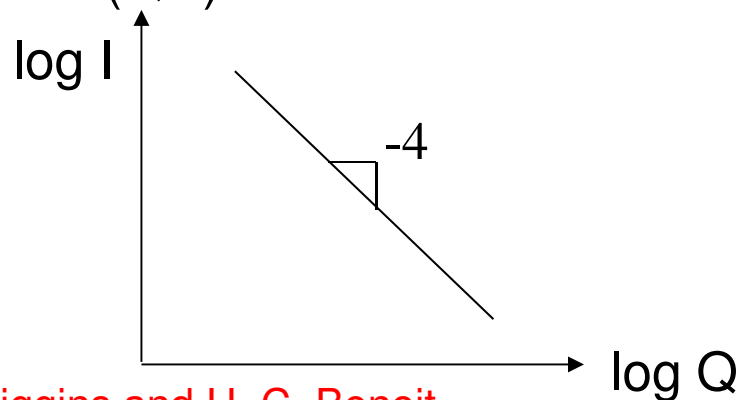
The form factor of a sphere (radius of R)

$$\left(\frac{d\sigma}{d\Omega}\right)_v = \phi_{\text{sphere}} V_{\text{sphere}} \Delta\rho^2 \frac{9}{(QR)^6} [\sin(QR) - QR \cdot \cos(QR)]^2$$

For  $QR \gg 1$ , i.e, focusing at smooth interface between two bulks 3-dimensionally

$$\left(\frac{d\sigma}{d\Omega}\right)_v \sim \frac{N_{\text{sphere}} V_{\text{sphere}}^2 \Delta\rho^2}{V} \frac{9}{2(QR)^4} = \frac{2\pi\Delta\rho^2 S_T}{VQ^4} \rightarrow \text{total surface area}$$

$$Q^4 \left(\frac{d\sigma}{d\Omega}\right)_v = \frac{2\pi\Delta\rho^2 S_T}{V}$$



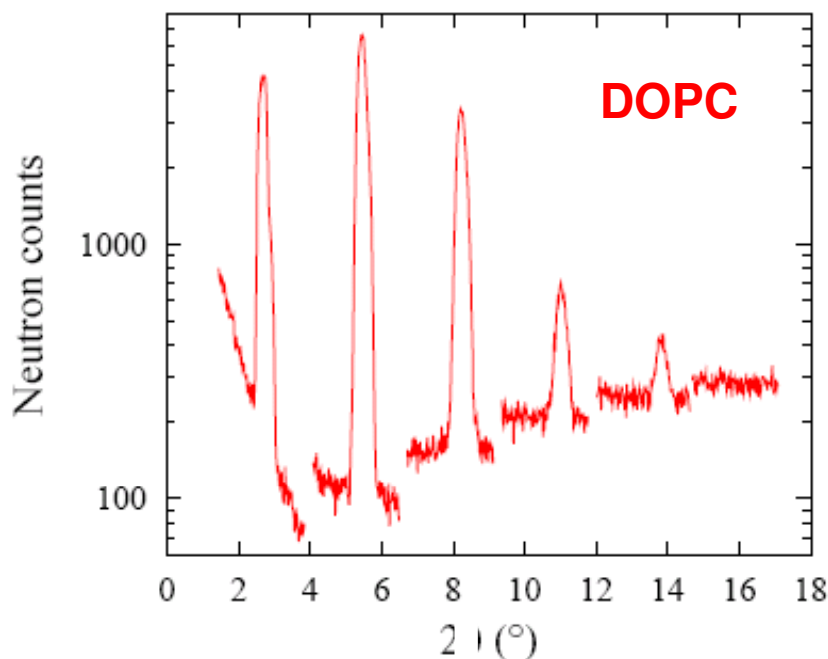
## SANS Analysis – Model Independent Scaling Method

If  $I_m(Q)$  scales with  $Q^{-n}$  over a wide range of  $Q$ . The “possible” structure can be obtained with some knowledge of the systems.

<b>particles</b>	<b>n</b>
•Long Rigid rod	1
•Smooth 2-D Objects	2
•Linear Gaussian Chain	2
•Chain with Excluded Volume	5/3
•Interfacial Scattering from 3-D Objects with Smooth Surface (Porod regime)	4
with fractal Surface	3 ~ 4

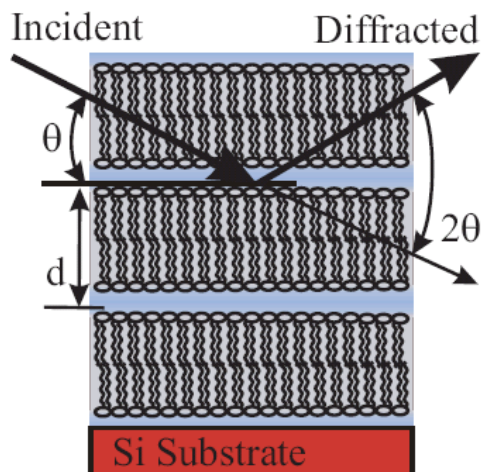
“Polymers and Neutron Scattering” J. S. Higgins and H. C. Benoit

# Small Angle Diffraction

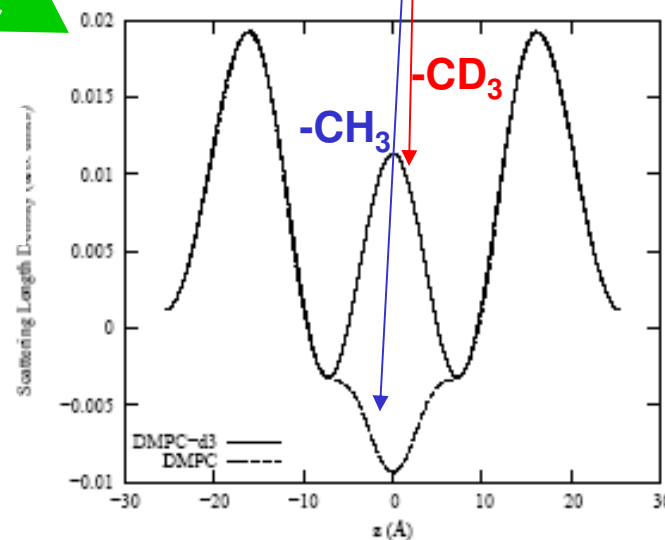


$$h \frac{Q}{2\pi} = \frac{h}{d} = \frac{2h \sin(\theta/2)}{\lambda}$$

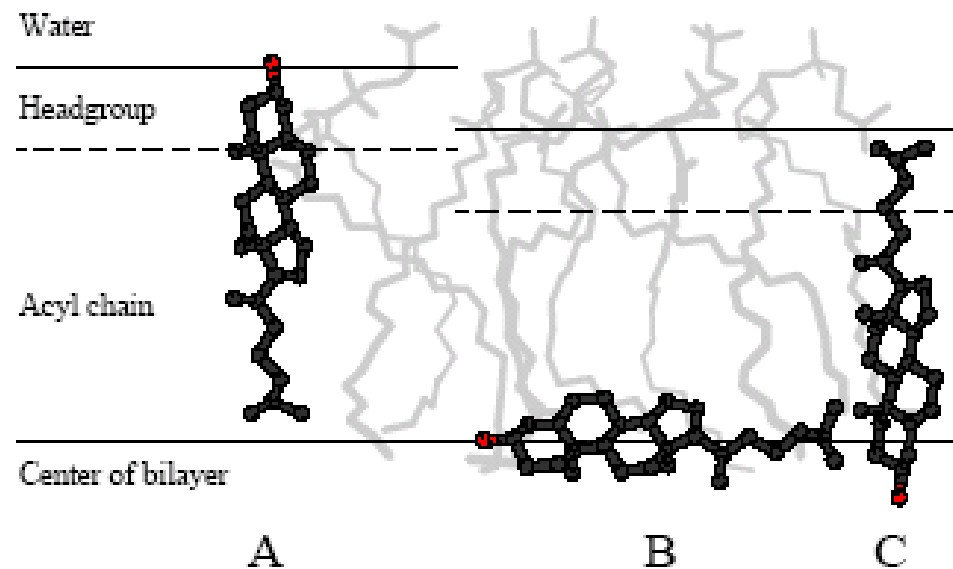
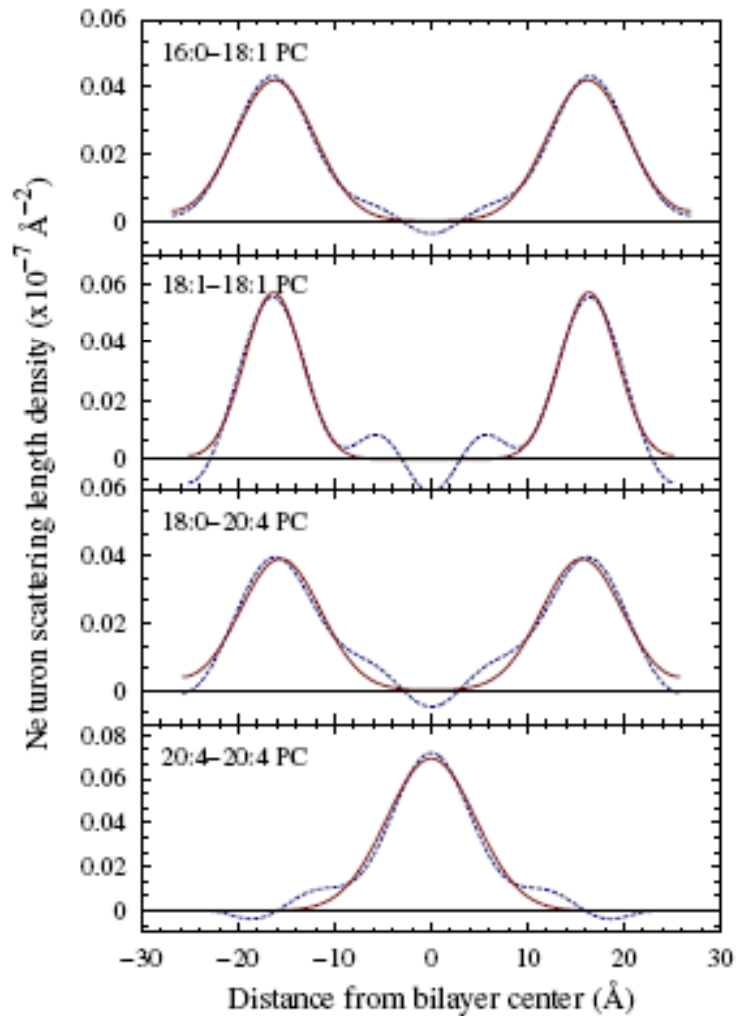
$$I \propto |F(Q)|^2 = \left| \sum b_i e^{-iQ \cdot z_i} \right|^2$$



$$\rho(z) = \rho_0 + \sum_{h=1}^{h_{max}} F_h \cos(2\pi h z / d)$$



## Example: Orientation of Cholesterol in Biomembrane



Harroun, T. A. et al. *Biochemistry* 45 (2006) 1227-1233.

## Take Home Messages

- SANS is a powerful tool to study the **global structures** of **isotropic systems** in the length scales ranging from **10 to 1000 Å**.
- It is easy to **vary the contrast** in a neutron scattering experiment **without significantly changing the chemical properties** of the systems.
- The scattering function is proportional to the product of **contrast factor**, **form factor** (intraparticle scattering function) and **structure factor** (interparticle interaction).
- Structural parameters can be obtained through **model dependent** approaches, while **model independent scaling law** can also reveal **possible morphology** of the studied system.